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Project **A1**

Electroweak Sudakov logarithms

New results for the form factor
in a massive $U(1)$ model and $U(1) \times U(1)$ with mass gap

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I Why logarithmic 2-loop results in EW theory?

Electroweak (EW) precision physics

- experimentally measured by now at energy scales up to $\sim M_{W,Z}$
- future generation of accelerators (LHC, LC) \rightarrow TeV region

Electroweak radiative corrections

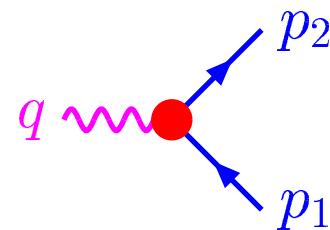
at high energies $\sqrt{s} \sim \text{TeV} \gg M_{W,Z}$

Kühn et al. '00, '01; Fadin et al. '00;
Denner et al. '01, '03; B.F. et al. '03;
Pozzorini '04

large negative corrections in exclusive cross sections

EW corrections dominated by Sudakov logarithms $\alpha^n \ln^{2n}(s/M_W^2)$

Form factor F of vector current:



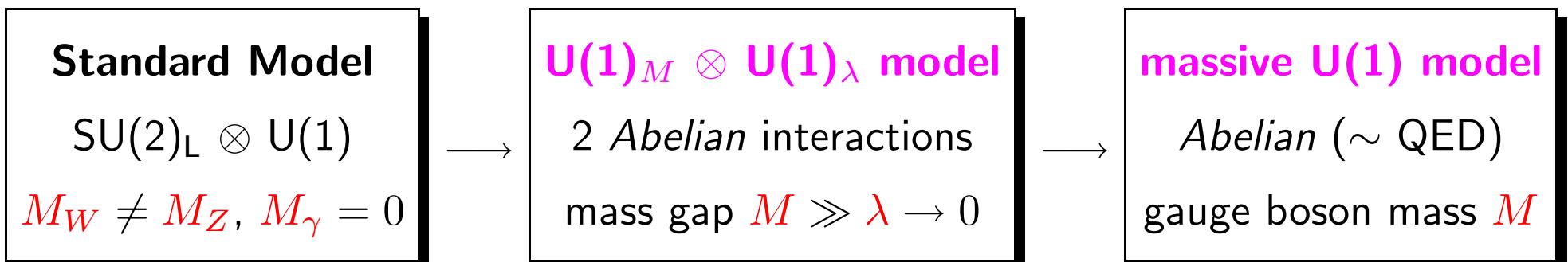
$$= \bar{u}(p_2) \gamma^\mu u(p_1) \cdot \mathbf{F} + \dots$$

High energy behaviour → *Sudakov limit*

- momentum transfer $|q^2| \equiv Q^2 \gg M^2$
- neglect fermion masses
- *logarithmic approximation*: neglect terms suppressed by a factor of M^2/Q^2
↪ works well for 2-loop n_f contribution where the exact result in M^2/Q^2 is known

B.F., Kühn, Moch '03

Simplified models



II Massive U(1) form factor

Form factor in perturbation theory: $F = 1 + \alpha F_1 + \alpha^2 F_2 + \dots$

large radiative corrections for $Q \sim \text{TeV} \rightarrow$ sum up large logarithms to all orders in α

Evolution equation in logarithmic approximation: Sen '81; Collins '89; Korchemsky '89; ...

$$\frac{\partial F(Q^2)}{\partial \ln Q^2} = \left[\int_{M^2}^{Q^2} \frac{dx}{x} \gamma(\alpha(x)) + \zeta(\alpha(Q^2)) + \xi(\alpha(M^2)) \right] F(Q^2)$$

Solution \rightarrow resummation (schematically):

$$F = 1 + \alpha (\ln^2 + \ln + \text{const}) + \alpha^2 (\ln^4 + \ln^3 + \ln^2 + \ln + \text{const}) + \dots$$

$$\rightarrow (1 + \alpha \cdot \text{const} + \alpha^2 \cdot \text{const} + \dots) \exp\left(\alpha (\ln^2 + \ln) + \alpha^2 (\ln^3 + \ln^2 + \ln) + \dots\right)$$

Anomalous dimensions γ, ζ, ξ from 1-loop calculation & massless 2-loop result

\Rightarrow NNLL approximation of F_2 known: $\alpha^2 (\ln^4 + \ln^3 + \ln^2)$

Massive U(1) form factor in 2-loop approximation

Known from resummation & full calculation of n_f contribution: $(n_f = \# \text{ fermions})$

$$\begin{aligned} \alpha^2 F_2 = \left(\frac{\alpha}{4\pi}\right)^2 & \left[+\frac{1}{2} \ln^4\left(\frac{Q^2}{M^2}\right) - \left(\frac{4}{9}n_f + 3\right) \ln^3\left(\frac{Q^2}{M^2}\right) \right. \\ & + \left(\frac{38}{9}n_f + \frac{2}{3}\pi^2 + 8\right) \ln^2\left(\frac{Q^2}{M^2}\right) \\ & - \left.\left(\frac{34}{3}n_f + \dots\right) \ln\left(\frac{Q^2}{M^2}\right) + \left(\frac{16}{27}\pi^2 + \frac{115}{9}\right)n_f + \dots\right] \end{aligned}$$

Kühn, Moch, Penin, Smirnov '01
B.F., Kühn, Moch '03

- growing coefficients with alternating sign:

$$\begin{aligned} & -0.4n_f \ln^3 + 4.2n_f \ln^2 - 11.3n_f \ln + 18.6n_f \\ & + 0.5\ln^4 - 3\ln^3 + 14.6\ln^2 - \dots \ln + \dots \end{aligned}$$

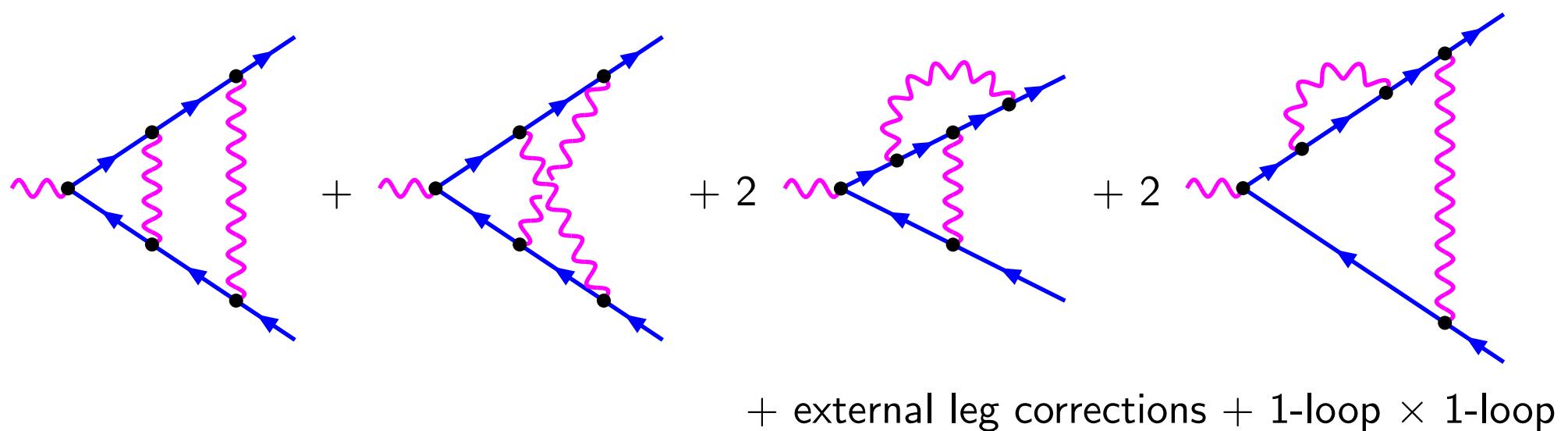
- $Q \sim 1 \text{ TeV} \rightarrow +\ln^4 \sim -\ln^3 \sim +\ln^2$

→ large cancellations between logarithmic terms

Complete 2-loop corrections in logarithmic approximation necessary.

Massive U(1) form factor in 2-loop approximation: calculation ($n_f = 0$)

- complete 2-loop result → loop calculation (*independent* of evolution equation)
- 2-loop vertex diagrams (massless fermions, massive bosons, 1 external scale):



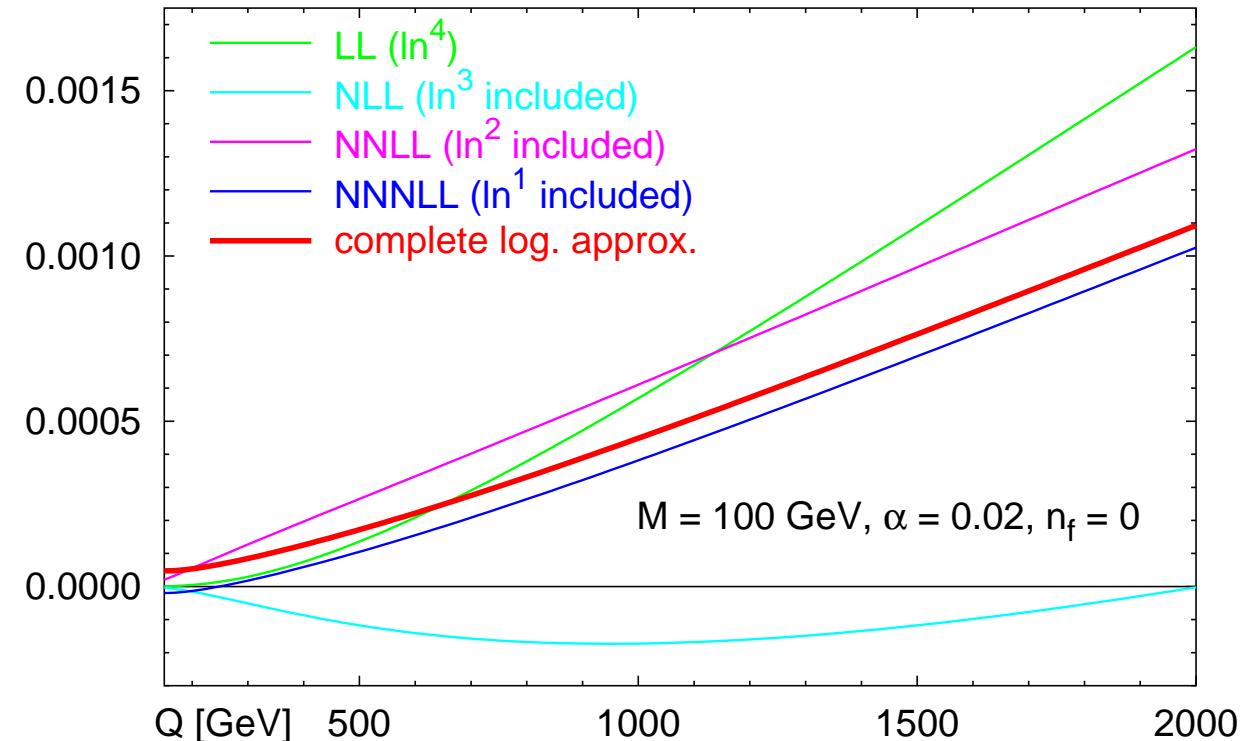
- evaluation of scalar diagrams: expansion by regions
- algorithms in FORM & Mathematica

Beneke, Smirnov '97

Massive U(1) form factor in 2-loop approximation: result ($n_f = 0$)

$$\alpha^2 F_2 = \left(\frac{\alpha}{4\pi}\right)^2 \left[\begin{aligned} & + \frac{1}{2} \ln^4 \left(\frac{Q^2}{M^2} \right) && \text{agreement ✓} \\ & - 3 \ln^3 \left(\frac{Q^2}{M^2} \right) \\ & + \left(\frac{2}{3} \pi^2 + 8 \right) \ln^2 \left(\frac{Q^2}{M^2} \right) \\ & - \left(-24\zeta_3 + 4\pi^2 + 9 \right) \ln \left(\frac{Q^2}{M^2} \right) \\ & + 256 \operatorname{Li}_4 \left(\frac{1}{2} \right) + \frac{32}{3} \ln^4 2 - \frac{32}{3} \pi^2 \ln^2 2 - \frac{52}{15} \pi^4 + 80\zeta_3 + \frac{52}{3} \pi^2 + \frac{25}{2} \end{aligned} \right]$$

new!



size of coefficients: $+0.5 \ln^4 - 3 \ln^3 + 14.6 \ln^2 - 19.6 \ln + 26.4$
at $Q = 1 \text{ TeV}$: $+225 \quad -293 \quad +309 \quad -90.4 \quad +26.4$

III $\text{U}(1) \times \text{U}(1)$ model with mass gap

EW theory: massive and massless gauge bosons

↪ consider $\text{U}(1)_M \times \text{U}(1)_\lambda$ model with 2 different masses $M \gg \lambda \rightarrow 0$

- pure $\text{U}(1)_M$: form factor $F(\alpha, Q, M)$
- pure $\text{U}(1)_\lambda$: form factor $F(\alpha', Q, \lambda)$
 - known from massive U(1) result ($M \rightarrow \lambda$, $\alpha \rightarrow \alpha'$)
 - IR singularities regularized by λ
- combined $\text{U}(1)_M \times \text{U}(1)_\lambda$: $\hat{F}(\alpha, \alpha', Q, M, \lambda)$ for $Q \gg M \gg \lambda \rightarrow 0$
 - **Factorization of IR singularities:**

$$\hat{F}(\alpha, \alpha', Q, M, \lambda) = \underbrace{F(\alpha', Q, \lambda)}_{\text{IR singular}} \underbrace{\tilde{F}(\alpha, \alpha', Q, M)}_{\text{IR finite}} + \mathcal{O}\left(\alpha \alpha' \frac{\lambda^2}{M^2}\right)$$

Factorization of $\mathbf{U}(1) \times \mathbf{U}(1)$ form factor: results ($n_f = 0$)

$$\hat{F}(\alpha, \alpha', Q, M, \lambda) = F(\alpha', Q, \lambda) \tilde{F}(\alpha, \alpha', Q, M) + \mathcal{O}\left(\alpha\alpha' \frac{\lambda^2}{M^2}\right)$$

$$\Rightarrow \tilde{F}(\alpha, \alpha', Q, M) = \lim_{\lambda \rightarrow 0} \frac{\hat{F}(\alpha, \alpha', Q, M, \lambda)}{F(\alpha', Q, \lambda)} = \lim_{\varepsilon \rightarrow 0} \frac{\hat{F}_\varepsilon(\alpha, \alpha', Q, M, 0)}{F_\varepsilon(\alpha', Q, 0)}$$

↪ set $\lambda = 0$ and calculate $\hat{F}_\varepsilon(\alpha, \alpha', Q, M, 0)$ in dimensional regularization

Calculation of 2-loop diagrams with 1 massive and 1 massless gauge boson →

$$\tilde{F}(\alpha, \alpha', Q, M) = F(\alpha, Q, M) \times$$

$$\left\{ 1 + \frac{\alpha\alpha'}{(4\pi)^2} \left[\underbrace{\left(48\zeta_3 - 4\pi^2 + 3 \right)}_{21.2} \ln\left(\frac{Q^2}{M^2}\right) + \underbrace{\frac{7}{45}\pi^4 - 84\zeta_3 + \frac{20}{3}\pi^2 - 2}_{-22.0} \right] \right\}$$

⇒ interference terms are finite ↽ IR singularities factorize

⇒ additional terms contain only single logarithm \ln^1

IV Summary & outlook

Massive U(1) form factor

- simple model with massive gauge bosons
- complete 2-loop result in logarithmic approximation ✓

U(1) \times U(1) model with mass gap

- first step towards EW theory with massive & massless gauge bosons
- factorization of IR singularities shown explicitly ✓

Outlook

- extend to non-Abelian models: SU(2), SU(N)
- consider effects from symmetry breaking: Higgs contributions, $M_Z > M_W$
- 4-fermion scattering amplitude
- full electroweak standard model